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## Reflections on Quantum Computing



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# Reflections on Quantum Computing 

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#### Abstract

In this rather speculative note three problems pertaining to the power and limits of quantum computing are posed and partially answered: a) when are quantum speedups possible?, b) is fixed-point computing a better model for quantum computing?, c) can quantum computing trespass the Turing barrier?


## 1 When are quantum speedups possible?

This section discusses the possibility that speedups in quantum computing can be achieved only for problems which have a few or even unique solutions [12]. For instance, this includes the computational complexity class UP [15]. Typical examples are Shor's quantum algorithm for prime factoring [18] and Grover's database search algorithm [13] for a single item satisfying a given condition in an unsorted database (see also Gruska [14]).

In quantum complexity, one popular class of problems is BQP, which is the set of decision problems that can be solved in polynomial time (on a quantum computer) so that the correct answer is obtained with probability at least $\frac{1}{2}$ on all instances. Both Shor's and Grover's problems are in BQP. The classical complexity of the primality problem is in NP $\cap$ co-NP and the unsorted database search problem is in P. However, Grover's quantum algorithm runs in time proportional to $\sqrt{n}$ for databases of size $n$, which is somewhat surprising since the classical lower bound is $\Omega(n)$ (for inputs of size $n$ the algorithm runs in time at least proportional to $n$ ).

Problems which can be efficiently solved by utilizing quantum parallelism may belong to a complexity class which we might call "quasi-UP". This class is characterized by a solution space which is small (say, polynomial) with respect to the dimension of the Hilbert space (the exponent of the number of qubits) involved. The class quasi-UP is very similar to the complexity class fewP $\subseteq$

[^0]NP that classifies problems with a fixed polynomial number of solutions per input size [15], so one can conjecture the following Quasi-UP-thesis:

The class fewP is a subset of the class BQP, which is a subset of quasi-UP.

We also suspect that neither NP nor BQP is a subset of the other. Here, BQP contains all of the bounded-error probabilistic polynomial-time problems (the class BPP), which potentially contains some co-NP problems not in NP. Also, most problems belonging to the class non-deterministic polynomial time NP are typically in another, dual regime: there, the number of conceivable solutions is large with respect to the number of bits involved to define the problem. Most NP-complete problems fall into this category, which is the primary reason we believe NP $\nsubseteq \mathbf{B Q P}$.

We now give one simple example of how the quasi-UP-thesis can be applied. Suppose we are interested in the intensively-studied traveling salesman problem (TSP) of finding the cheapest trip in cost through all of the nodes of a map. Suppose we know (as naturally suspected) that our problem instances (i.e., in the "real-world") have a unique best solution (or, at worst, a few equal optimal solutions). For these types of inputs, our restricted TSP problem is in few $\mathbf{P}^{1}$ and thus we can expect to have an efficient quantum search algorithm.

Although we cannot give a direct proof of this quasi-UP-thesis, some informal arguments can be brought forward in its support. Efficient algorithms in quantum computing make use of the quantum parallelism. Yet in order to be able to extract a classically useful solution from the resulting quantum state, one has to extract the information by proper phase transformations and interference. And it is interference - the buildup of phases at points which indicate the problem solutions-we are mostly concerned about. Interference guarantees that the result of the quantum calculation can be effectively read out of the superposition of states.

A problem allowing only a single solution therefore has a better chance to be solvable by a quantum algorithm. In particular, in the interference phase the single solution could allow for a higher contrast ${ }^{2}$ and thus a better detection efficiency than a situation which would allow for many solutions. Thus, for example, suppose we apply a quantum algorithm to an unsorted database problem that allows up to a fixed constant number of matches. Then our success with a quantum algorithm will probably deteriorate even though the probability of randomly finding a match increases.

It therefore appears to be not totally unreasonable to speculate that detection efficiency might drop linearly with the number of solutions, as for a problem with $n$ items the contrast drops like $O\left(1 / e^{n}\right)$. This could make most NP problems effectively intractable for quantum algorithms. Thus the ability to enhance contrast and detector efficiency for problems in NP appears to be

[^1]one of the most crucial steps a quantum algorithm has to cope with. The task here is formidable-to single out the most favorable solution out of a sea of possible but non-optimal ones.

## 2 Fixed-point quantum computations

One radically new possibility for quantum computations would be an approach based on fixed-point computations. Stated differently, a quantum computation may arrive at results which are fixed-points of a unitary operator or, more precisely, eigenstates of unitary or Hermitian operators with eigenvalue 1.

One task which is impossible within the domain of classical computation but almost trivial for quantum algorithms is the solution of diagonalization problems. Diagonalization is based on bit switches from true to false and vice versa. In order to solve this problem for a single qubit it can be implemented by a unitary and self-adjoint not-operator

$$
\widehat{D}=n o t=\left(\begin{array}{ll}
0 & 1 \\
1 & 0
\end{array}\right)
$$

The eigenstates of $\widehat{D}$ are

$$
\begin{equation*}
|I\rangle,|I I\rangle=\frac{1}{\sqrt{2}}\left[\binom{1}{0} \pm\binom{ 0}{1}\right] \tag{1}
\end{equation*}
$$

with the eigenvalues +1 and -1 , respectively. The solution of the diagonalization problem is the eigenvector $|I\rangle$ associated with the eigenvalue 1. This is a mixture or superposition of the classical bit states true and false.

Other, more general problems may be solvable by applying the fixed-point theorem of computability theory to quantum computations. The strategy is to assume a normal operator $A$ which, due to some yet unknown procedure, encodes an algorithm $\varphi_{A}$. Thereby, we seek an unknown eigenstate $a_{1}$ of $A$ with eigenvalue 1. Let us assume that we start with a totally "unbiased" state 1 which, in matrix notation, is just represented by the unit matrix. Indeed, if we have information about the solution we may prepare the state in such a way that the outcome of the fixed-point is more likely. (In the extreme case we know the solution before the measurement and prepare the initial state to be exactly the fixed-point state. The outcome of the fixed-point state is thus certain.) As $A$ is measured, $a_{1}$ is obtained if we observe the eigenvalue 1 . We may now identify the fixed-point solution $a_{1}$ with the algorithm $\varphi_{A}$.

## 3 Computing the uncomputable?

One fundamental result of theoretical computer science is Turing's proof (in [20]) that it is undecidable to determine whether a general computer program will halt or not. This is formally known as the halting problem. We can restrict our
attention to Turing machines, since they are equivalent in computational power to any "conventional" computer $[3,1]$. In what follows we present an attempt to trespass the Turing barrier. The method discussed might in principle allow us to "solve" the halting problem (for another proposal see Mitchison and Josza, [16]). Thereby we are well aware of the fact that for all practical purposes (Bell [2]) this goal will remain unreachable, at least within quantum computing.

Assume that it is possible to design a halting qubit which indicates whether a computation has actually reached a state associated with a halting condition. Assume further that the halting qubit starts in its non-halting state and, since the evolution is unitary, the buildup of the amplitude is continuous in time.

In such a case, the halting qubit acquires a halting component which is non-zero even in finite time. Therefore, a detection of a halting computation at small time scales is conceivable even if the associated classical computation lasts "very" long. The price to be paid is the "very small" amplitude and, associated with it, a correspondingly small chance of detection.

To be a little bit more precise we need some rudiments of algorithmic information theory (see Chaitin [6, 7], Calude [4]). We will work with programs with no input which produce binary strings as outputs. For any $n$ we denote by $P_{n}$ a program of length $n$ that halts and produces the longest string among all outputs produced by all programs of length $n$ that eventually stop. We denote by $\Sigma(n)$ the length of the output produced by $P_{n}$. Here $\Sigma$ is the busy beaver function [19, 8]: it grows faster than every computable function of $n$. Let $H$ be the program-size complexity, that is the length of the smallest universal program generating a particular binary string.

Assume that any program which halts requires a running time at least proportional to the length of its output. If an $n$-bit program $p$ halts, then the time $t$ it takes to halt satisfies $H(t) \leq n+c$. So if $p$ has run for time $T$ without halting, and $T$ has the property that if $t \geq T$, then $H(t)>n+c$, then $p$ will never halt. This shows that the running times of the programs in the sequence $P_{1}, P_{2}, \ldots P_{n}, \ldots$ grow faster than any computable function.

We are now ready to present the argument. Let us assume the halting qubit is represented by

$$
|H a l t\rangle=c_{h}(t)|h\rangle+c_{n}(t)|n\rangle,
$$

where $|h\rangle,|n\rangle$ represent the halting state and non-halting state and $c_{h}(t), c_{n}(t)$ are time-dependent amplitudes thereof, respectively.

Initially, let $\left|c_{h}(t)\right|=\left|c_{n}(t)\right|-1=0$. As a worst-case scenario derived from the above analysis, for a linear buildup of the amplitude we obtain

$$
\left|c_{h}(t)\right|^{2} \propto(\Sigma(H(n)+O(1)))^{-1}
$$

The setup of a detection of $|H a l t\rangle$ is a simple transmission measurement of the halting qubit. Although the buildup may be very slow, there is a nonvanishing chance to obtain a solution of the halting problem in finite time. Of course, the solution is probabilistic (one can argue that all mathematical proofs or computer programs are ultimately probabilistic, see Davis [9], De Millo, Lipton, Perlis [11]), but goes beyond the capability of any classical computation:
even the best probabilistic algorithms are not able to achieve this computational power (by a classical result [10], probabilistic algorithms are equivalent to Turing machines).

Let us finally notice that by virtue of the same information-theoretic argument, the possibility of time-travel (see, for example, Nahin [17]) would not solve the halting problem, unless one could travel back and forth in time at a pace exceeding the growth of any computable function.

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[^1]:    ${ }^{1}$ Our problem is probably not NP-hard.
    ${ }^{2}$ The number of solutions seem to influence the relative minimal and maximal particle intensities and the width thereof, as they are observed in an interference pattern depicting the average number of clicks in a detector measuring the output of a quantum computer.

