Boole-Bell-Type Inequalities in Mathematica

Stefan Filipp and Karl Svozil

Institut für Theoretische Physik, Technische Universität Wien Wiedner Hauptstraße 8-10/136, A-1040 Vienna, Austria e-mail: svozil@tuwien.ac.at

Abstract

Classical Pitowsky correlation polytopes with particular emphasis on the Minkowski-Weyl representation theorem are reviewed. A numerical study of the generation of optimal Boole-Bell type inequalities for arbitrary experimental setups is presented. These inequalities represent the faces of classical Pitowsky correlation polytopes [1, 2], which are based on the "conditions of possible experience" stated by George Boole in the 19th century[3]. We introduce CddIF, a Mathematica package created as an interface between Mathematica and the cdd program by Komei Fukuda[4], which represents a highly efficient method to solve the hull problem for general classical correlation polytopes.

Contents

1	Bool	e-Bell Type Inequalities	3
	1.1	Simple urn model	3
	1.2	Geometrical interpretation	3
	1.3	Minkowski-Weyl representation theorem	4
	1.4	From vertices to inequalities	4
	1.5	From inequalities to vertices	5
	1.6	Quantum mechanical context	5
2	Insta	ıllation	7
	2.1	Mathematica	7
	2.2	cdd	8
	2.3	Installation on windows-platforms	8
3	Desc	ription Of Functions	9
	3.1	CddFormat	9
	3.2	ToCddExtFile	9
	3.3	TruthTable	10
	3.4	RunCdd	11
	3.5	ShowVRep	12
	3.6	ConvToHRep	12
	3.7	ReadInHRep	13
	3.8	GetInequFromHRep	13
	3.9	InequToRead	14
	3.10	GetViolInequalities	14
	3.11	PlotInequalities	15
		ContPlotInequalities	17
	3.13	Cdd	19
4	Exar	nples	20
	4.1	Three particles and two measurement directions	20
	4.2	Violations of inequalities	21
	4.3	Graphical representation	22
	4.4	Two particles and three measurement directions	24
	4.5	Violations of inequalities	24
	16	Graphical representation	26

1 Boole-Bell Type Inequalities And Their Geometric Representation

In the middle of the 19th century the English mathematician George Boole formulated a theory of "conditions of possible experience" [5, 3, 6, 1, 7]. These conditions are related to relative frequencies of logically connected events and are expressed by certain equations or inequalities. More recently, similar equations for a particular setup which are relevant in the quantum mechanical context have been discussed by Clauser and Horne and others [8, 9, 10]. Pitowsky has given a geometrical interpretation in terms of correlation polytopes [11, 1, 2, 7].

1.1 Simple urn model

Consider an urn containing some balls of different colors and styles: Each ball can be described by two properties, let us say "yellow" and "wooden", so each ball can have the property "yellow" or the property "wooden", but it can also have both - "yellow and wooden". The state of the urn can be given by the probabilities to draw a ball with one of these properties: p_1 is the proportion of yellow balls in the urn, p_2 the proportion of wooden ones and p_{12} denotes the proportion of yellow and wooden balls. If there are enough balls in the urn these proportions are in fact the probabilities to get a ball with the special property by drawing. Clearly the inequalities

$$0 \le p_{12} \le p_2 \le 1$$
 and $0 \le p_{12} \le p_1 \le 1$ (1)

are fulfilled by the proportions and so p_1 , p_2 and p_{12} can be seen as probabilities of two events and their joint event only if these inequalities are satisfied. Simply by taking some appropriate numbers ($p_1 = 0.6$, $p_2 = 0.72$ and $p_{12} = 0.32$) we can see, that equations (1) are not sufficient. If we take the probability to draw a ball which is either yellow or wooden ($p_1 + p_2 - p_{12}$) into consideration, a new inequality can be found that is not satisfied by the numbers chosen:

$$0 \le p_1 + p_2 - p_{12} \le 1 \tag{2}$$

It can be shown that the inequalities (1) and (2) are necessary and sufficient for the numbers p_1 , p_2 and p_{12} to represent probabilities of two events and their joint [1].

1.2 Geometrical interpretation

Itamar Pitowsky [11, 1, 2, 7] has suggested a geometric interpretation. Let us take the set of all numbers (p_1, p_2, p_{12}) satisfying the inequalities stated above as a set of vectors in a three-dimensional real space. This procedure yields a closed convex polytope with vertices (0,0,0), (1,0,0), (0,1,0) and (1,1,1) (cf. Figure 1). The extreme points (vertices) can be interpreted as follows:

- (0,0,0) is a case where no yellow and no wooden balls are in the urn at all,
- (1,0,0) is representing the configuration that all balls are yellow and no one is wooden.
- (0,1,0) is representing the configuration that all balls are wooden and no one is yellow.
- (1,1,1) is a case with only yellow and at the same time wooden balls.

Figure 1: Polytope associtated with the urn model

1.3 Minkowski-Weyl representation theorem

The Minkowski-Weyl representation theorem (e.g., [12, p. 29]) states that compact convex sets are "spanned" by their extreme points; and furthermore that the representation of this polytope by the inequalities corresponding to the planes of their faces is an equivalent one.

Stated differently, every convex polytope in an Euclidean space has a dual description: either as the convex hull of its vertices (V-representation), or as the intersection of a finite number of half-spaces, each one given by a linear inequality (H-representation) This equivalence is known as the *Weyl-Minkowski* theorem.

The problem to obtain all inequalities from the vertices of a convex polytope is known as the *hull problem*. One solution strategy is the Double Description Method [13] which we shall use but not review here.

1.4 From vertices to inequalities

For the above simple urn model, the inequalities are rather intuitive and can be easily obtained by guessing. This is impossible in the general case involving more events and more joint probabilities thereof. In order to obtain the relevant inequalities—Boole's "conditions of possible experience"—we have to find a hopefully constructive way to derive them.

Recall that a vector is an element of the polytope if and only if it can be represented as a certain bounded convex combination, i.e., a bounded linear span, of the vertices. More precisely, let us denote the *convex hull* conv(K) of a finite set of points $K = \{x_1, ..., x_n\} \in R^d$ by

$$\operatorname{conv}(K) = \left\{ \lambda_1 \mathbf{x_i} + \dots + \lambda_n \mathbf{x_n} \mid n \ge 1, \lambda_i \ge 0, \sum_{i=1}^n \lambda_i = 1 \right\}. \tag{3}$$

In the probabilistic context, the coefficients λ_i are interpreted as the probability that the event represented by the extreme point $\mathbf{x_i}$ occurs, whereby K represents the complete set of all atoms of a Boolean algebra. The geometric interpretation of K is the set of all extreme points of the correlation polytope.

In summary, the connection between the convex hull of the extreme points of a correlation polytope and the inequalities representing its faces is guaranteed by the Minkowski-Weyl representation theorem. A constructive solution of the corresponding hull problem exists (but is NP-hard [2]).

For the special urn model introduced above this means that any three numbers (p_1 , p_2 and p_{12}) must fulfill an equation dictated by Kolmogorov's probability axioms [14]:

$$(p_1, p_2, p_{12}) = \lambda_1(0, 0, 0) + \lambda_2(0, 1, 0) + \lambda_3(1, 0, 0) + \lambda_4(1, 1, 1) = (\lambda_2 + \lambda_4, \lambda_3 + \lambda_4, \lambda_4). \tag{4}$$

It is important here to realize that these logical possibilities are exhaustive, That is, by definition, there cannot be any other classical case which is not already included in the above possibilities (0,0,0),(1,0,0),(0,1,0),(1,1,1). the third "component bit" of the vector is a function of the first components. Actually, the function is a multiplication, since we are dealing with the classical logical "and" operation here. (If one or more cases are omitted, the corresponding polytope would not be optimal; i.e., it would be embedded in the optimal one.) Therefore, any "state" of a physical system represented by a probability distribution has to satisfy the constraint

$$\lambda_1 + \lambda_2 + \lambda_3 + \lambda_4 = 1. \tag{5}$$

The extreme cases $\lambda_i = 1, \lambda_j = 0$ for $i \in \{1, 2, 3, 4\}$ and $j \neq i$ just correspond to the vertices spanning the classical correlation polytope as the convex sum (3).

A generalization to arbitrary configurations is straightforward. To solve the hull problem for more general cases, an efficient algorithm has to be used. There are some algorithms to solve this problem, but they run in exponential time in the number of events, thus it can be solved only for small enough cases to get a solution in conceivable time.

1.5 From inequalities to vertices

Conversely, a vector is an element of the convex polytope if and only if its coordinates satisfy a set of linear inequalities which represent the supporting hyper-planes of that polytope. The problem to find the extreme points (vertices) of the polytope from the set of inequalities is dual to the hull problem considered above.

1.6 Quantum mechanical context

In the quantum mechanical case the elementary irreducible events are clicks in particle detectors and the probabilities have to be calculated using the formalism of quantum mechanics.

Figure 2: Experimental setting to test the violation of Boole - Bell type inequalities

As an example we take a source that produces pairs of spin- $\frac{1}{2}$ particles in a singlet-state $(|\psi\rangle = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle))$. The particles fly apart along the z axis and after the particles have separated, measurements on spin components along one out of two directions are made. If, for simplicity, the measurements are made in the x-y plane perpendicular to the trajectory of the particles, the direction of the measurement can be given by angles measured from the vertical x axis (α_1 and α_2 on the one side, β_1 and β_2 on the other side). On each side the measurement angle is chosen randomly for each pair of incoming particles and each measurement can yield two results - in $\frac{\hbar}{2}$ units: "+1" for spin up and "-1" for spin down (cf. Figure 2).

Deploying this configuration we get probabilities to find a particle measured along the axis specified by the angles α_1 , α_2 , β_1 and β_2 either in spin up or in spin down state - denoted as p_{a1} , p_{a2} , p_{b1} , p_{b2} - and we also take the joint event of finding a particle on one side at the angle α_1 (α_2) in a specific spin state and the other particle on the other side along the vector β_1 (β_2) in a specific spin state, denoted as p_{a1b1} , p_{a2b1} , p_{a1b2} and p_{a2b2} . To construct the convex polytope to this experiment we build up a truth table of all possible events using a "1" as "spin up is detected along the specific axis" and a "0" as "spin down is detected along the specific axis" (table 1). The rows of this table

α_1	α_2	β_1	β_1	$\alpha_1\beta_1$	$\alpha_1\beta_2$	$\alpha_2\beta_1$	$\alpha_2\beta_2$
0	0	0	0	0	0	0	0
1	0	0	0	0	0	0	0
0	1	0	0	0	0	0	0
1	1	0	0	0	0	0	0
0	0	1	0	0	0	0	0
1	0	1	0	1	0	0	0
0	1	1	0	0	0	1	0
1	1	1	0	1	0	1	0
0	0	0	1	0	0	0	0
1	0	0	1	0	1	0	0
0	1	0	1	0	0	0	1
1	1	0	1	0	1	0	1
0	0	1	1	0	0	0	0
1	0	1	1	1	1	0	0
0	1	1	1	0	0	1	1
1	1	1	1	1	1	1	1

Table 1: Truth table for four propositions

are then identified with the vertices of the convex polytope. By using the Minkowski-Weyl theorem and by solving the hull problem, the vertices determine the hyper-planes confining the polytope, i.e. the inequalities which the probabilities have to satisfy. As a result the following inequalities are gained:

$$0 \le p_{aibi} \le p_{ai} \le 1, 0 \le p_{aibi} \le p_{bi} \le 1 \quad i = 1, 2
p_{ai} + p_{bi} - p_{aibi} \le 1 \quad i = 1, 2$$
(6)

$$-1 \le p_{a1b1} + p_{a1b2} + p_{a2b2} - p_{a2b1} - p_{a1} - p_{b2} \le 0
-1 \le p_{a2b1} + p_{a2b2} + p_{a1b2} - p_{a1b1} - p_{a2} - p_{b2} \le 0
-1 \le p_{a1b2} + p_{a1b1} + p_{a2b1} - p_{a2b2} - p_{a1} - p_{b1} \le 0
-1 \le p_{a2b2} + p_{a2b1} + p_{a1b1} - p_{a1b2} - p_{a2} - p_{b1} \le 0$$
(7)

The last four inequalities are known as *Clauser-Horne inequalities*. As noticed above the probabilities have to be seen in a quantum mechanical context. If we consider the singlet state of spin- $\frac{1}{2}$ particles $|\psi\rangle = \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$ it is well known that the probability to find the particles both either in spin up or in spin down states is given

by $P^{\uparrow\uparrow}(\theta)=P^{\downarrow\downarrow}(\theta)=\frac{1}{2}sin^2(\theta/2)$ - where θ is the angle between the measurement directions. The single event probability is clearly $p_i=\frac{1}{2}$. By choosing

$$a_1 = -\frac{\pi}{3}$$
 $a_2 = b_1 = \frac{\pi}{3}$ $b_2 = \frac{\pi}{3}$ (8)

as measurement directions, we get for $p = (p_{a1}, p_{a2}, p_{b1}, p_{b2}, p_{a1b1}, p_{a2b1}, p_{a1b2}, p_{a2b2})$:

$$p = (\frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{1}{2}, \frac{3}{8}, \frac{3}{8}, 0, \frac{3}{8}) \tag{9}$$

and one of the inequalities found in (7) is violated:

$$p_{a1b1} + p_{a1b2} + p_{a2b2} - p_{a2b1} - p_{a1} - p_{b2} = \frac{3}{8} + \frac{3}{8} + \frac{3}{8} - 0 - \frac{1}{2} - \frac{1}{2} = \frac{1}{8} > 0$$
 (10)

The generalization is straightforward. Violations of certain inequalities involving classical probabilisties—Boole's "conditions of possible experience" [3]—also appear in higher dimensions in configurations containing more particles and/or more measurement directions. We shall consider more examples below.

2 Installation

2.1 Mathematica

All functions described in the following section can be found in the *Mathematica*-package *cddif.m*. In general this package has to be loaded into the current *Mathematica*-kernel by the command <<'path to cddif.m'/cddif.m, short description and usage of the functions is available by entering ?'<function>'.

To guarantee a proper run of all functions it is necessary (and hopefully sufficient) that cdd is located in any directory listed in the PATH-variable (usually /bin, /usr/bin, /usr/local/bin, ...)¹ or in the current working directory, which can be shown by evaluating <code>Directory[]</code> or changed using the function <code>SetDirectory[directory_String]</code>. If you like to avoid the frequent use of this function you can append this command to the package-file <code>cddif.m</code> before the line <code>End[]</code> so that on each loading of the package the directory is set automatically to your personal working directory.

2.2 cdd

cdd is a C++ (ANSI C) implementation of the Double Description Method [13] for generating all vertices (i.e. extreme points) and extreme rays of a general convex polyhedron given by a system of linear inequalities and the dual problem of generating a system of linear inequalities given all vertices by Komei Fukuda[4].

At this point I only refer to the documentation of the program for the installation of the cdd - package, in particular to the file

¹ for setting environment variables look at the manual of set or env

cdd.readme included in the package and to the online documentation at http://www.ifor.math.ethz.ch/~afukuda/cddman/cddman.html. cdd is available for free and you can download it from the homepage of Komei Fukuda[4] (http://www.ifor.math.ethz.ch/~fukuda/cdd_home/cdd.html), where you can also find a manual to the usage of cdd, especially descriptions to the format of the inputand output-files and to options that can be passed to cdd.

2.3 Installation on windows-platforms

cdd is available to compile on UNIX-like systems, to run cdd on Windows - systems you can either try to compile it in Windows² or look somewhere else, if anybody has already done this work. Currently a version runnable on Windows platforms of cdd can be downloaded from http://tph.tuwien.ac.at/~svozil/cdd/cdd.exe. (a different compilation is at http://www.wis.kuleuven.ac.be/wis/algebra/kathleen/files/cdd061.exe). On Windows-systems *Mathematica* must be able to find cdd in a directory listed in the PATH-variable or in the current working directory, too. To set the PATH-variable in Windows 2000/NT go to the "control panel" and click on the "system properties", then click "advanced" and there is a place where the variable PATH is specified. Here you can add your path to cdd (separated by a semicolon). In WindowsME you need to go execute "msconfig" to get to the System Configuration Utility - in "Environment" you can edit the PATH-variable and for Windows98/95 you must edit the file "autoexec.bat" to get the path set.

Finally it can be necessary to rename the cdd - executable file (e.g. from *cdd061.exe*) to *cdd.exe*. The CddIF-Package uses *cdd* as default command to run cdd, using the function SetCddCmd [*cmd_String*] you can change this behavior. Like already stated above you can also add this commandline to the package-file *cddif.m* just before the End[]-statement to change the default command automatically when loading the package.

3 Description Of Functions

In this section all functions of the *CddIF* - package are listed. For each function the syntax including the necessary parameters (if parameters are optional, it has an "opt:" as prefix), a description and an example is given.

3.1 CddFormat

CddFormat [vertices, opt: options]

vertices (List): List of m vertices in n dimensions of the form

$$\{\{x_{11}, x_{12}, \ldots, x_{1n}\}, \{x_{21}, \ldots, x_{2n}\}, \ldots, \{x_{m1}, x_{m2}, \ldots, x_{mn}\}\}$$

²the manual says: "The program cdd.c is written in ANSI C, and thus it should run on personal computers without any changes if one uses a compiler supporting ANSI standard."

```
options (List): Options to cdd (e.g. adjacency, nondegenerate, minindex,...) - see documentation to cdd (http://www.ifor.math.ethz.ch/ fukuda/cddman/cddman.html)
```

Description: A list of vertices, which can be determined for example by TruthTable [..., *IncludeVars*→*False*], are converted to a format recognized by cdd. Additionally options to cdd can be declared.

Example:

```
In[]= CddFormat [ \{\{1,0,0\},\{0,1,0\},\{1,1,1\}\}\}]
Out[]= \{V-representation begin,\{3, 4, integer\},\{1, 1, 0, 0\},\{1, 0, 1, 0\},\{1, 1, 1, 1\},end\}
```

3.2 ToCddExtFile

```
ToCddExtFile [ file, vertices, opt: options ] or
ToCddExtFile [ file, particles, measurements, opt: options ]
```

file (String): Filename for output of H-representation (".ext"-suffix is automatically appended)

vertices (List): List of m vertices in n dimensions of the form

```
\{\{x_{11}, x_{12}, \ldots, x_{1n}\}, \{x_{21}, \ldots, x_{2n}\}, \ldots, \{x_{m1}, x_{m2}, \ldots, x_{mn}\}\}
```

options (List): Options to cdd (e.g. adjacency, nondegenerate, minindex,...) - see documentation to cdd (http://www.ifor.math.ethz.ch/ fukuda/cddman/cddman.html)

particles (Integer): Number of particles

measurements (Integer): Number of possible measurements to each particle (equivalent to number of detection angles)

Description: Creates a file with ".ext"-extension that contains the data of the given configuration to use in cdd. Eighter a list of vertices of the considered correlation polytop or the number of particles used and the possible measurements to each can be handed over. In the latter case the list of vertices is generated automatically.

Example:

```
In[] = ToCddExtFile [ "test",2,3 ]
```

Out[]= test.ext

writes the file "test.ext" to the current working directory, containing the vertices of the 2-particles 3-measurements configuration.

3.3 TruthTable

TruthTable [particles, measurements, opt: options]

particles (Integer): Number of particles

measurements (Integer): Number of possible measurements to each particle (equivalent to number of detection angles)

opt: options: The only possible option is IncludeVars → True/False. If IncludeVars → True is defined, the function includes a list of variables belonging to the given configuration as titles of the columns and output will be made in MatrixForm, otherwise a list containing all vertices is returned. Default is IncludeVars → True.

Description: Creates a truth table of the given configuration, containing all vertices of the corresponding correlation polytopes. For generating this table all possible single events are rated eighter 0 or 1 (i. e. true or false) and the joint events are evaluated using the logical *AND* operation.

Example:

In[]= TruthTable [2,2]
Out[]=

a1	a2	b1	b2	a1b1	a1b2	a2b1	a2b2
0	0	0	0	0	0	0	0
1	0	0	0	0	0	0	0
0	1	0	0	0	0	0	0
1	1	0	0	0	0	0	0
0	0	1	0	0	0	0	0
1	0	1	0	1	0	0	0
0	1	1	0	0	0	1	0
1	1	1	0	1	0	1	0
0	0	0	1	0	0	0	0
1	0	0	1	0	1	0	0
0	1	0	1	0	0	0	1
1	1	0	1	0	1	0	1
0	0	1	1	0	0	0	0
1	0	1	1	1	1	0	0
0	1	1	1	0	0	1	1
1	1	1	1	1	1	1	1

3.4 RunCdd

RunCdd [file]

file (String): File handed over to cdd as command parameter (automatically extended with "ext."-suffix.

Description: RunCdd [...] executes cdd (with file as parameter) and returns the corresponding H-representation, which can be used in various other functions like PlotInequalities [...] or GetViolInequalities [...].

Using this function you have to pay attention to the potentially long runtime in calculating the faces (i.e. the inequalities) of the correlation polytope. It can be more beneficial to use ToCddExtFile [file,...] to create a ".ext"-file, followed by executing cdd outside of *Mathematica* (eventually on a faster computer) to convert the date to H-representation stored in an ".ine"-file. Afterwards you can read in this file utilizing ReadInHRep [file]

Example:

```
In[]= RunCdd [ "test"]
Out[]= {{H-representation}, {begin}, {684,16,real}, {2,0,-
2,....},...., {end}},
    whereas 2-particles 3-measurement configuration is taken into consideration here.
```

3.5 ShowVRep

```
ShowVRep [ particles, measurements ] particles (Integer): Number of particles
```

measurements (Integer): Number of possible measurements to each particle (equivalent to number of detection angles)

Description: Shows the V-representation of a given configuration.

Example:

```
In[]= ShowVRep[ 2,2 ]
Out[]= {V-representation
    begin,
    {16, 9, integer},
    {1, 0, 0, 0, 0, 0, 0, 0, 0},{1, 1, 0, 0, 0, 0, 0, 0,
        0},
    {1, 0, 1, 0, 0, 0, 0, 0, 0},{1, 1, 1, 0, 0, 0, 0, 0,
        0},
    {1, 0, 0, 1, 0, 0, 0, 0, 0},{1, 1, 0, 1, 0, 1, 0, 0,
        0},
    {1, 0, 1, 1, 0, 0, 0, 1, 0},{1, 1, 1, 1, 1, 0, 1, 0, 1,
        0},
    {1, 0, 0, 0, 1, 0, 0, 0, 0},{1, 1, 1, 0, 0, 1, 0, 1,
        0},
    {1, 0, 0, 0, 1, 0, 0, 0, 0},{1, 1, 0, 0, 1, 0, 1, 0,
        0},
```

```
{1, 0, 1, 0, 1, 0, 0, 0, 1},{1, 1, 1, 0, 1, 0, 1, 0, 1}, {1, 0, 0, 0, 0, 0}, {1, 1, 0, 1, 1, 1, 1, 0, 0}, {1, 0, 1, 1, 1, 0, 0, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1}, end}
```

3.6 ConvToHRep

ConvToHRep [particles, measurements, opt: file, opt: options] or ConvToHRep [vertices, opt: file, opt: options]

particles (Integer): Number of particles

measurements (Integer): Number of possible measurements to each particle (equivalent to number of detection angles)

vertices (List): List of m vertices in n dimensions of the form

$$\{\{x_{11}, x_{12}, \ldots, x_{1n}\}, \{x_{21}, \ldots, x_{2n}\}, \ldots, \{x_{m1}, x_{m2}, \ldots, x_{mn}\}\}$$

file (String): Filename that is used for the conversion from a ".ext"-file to a ".ine"-file which is equivalent to a conversion from V-representation to H-representation. Default is "tmp".

options (List): Options to cdd (e. g. adjacency, nondegenerate, minindex,...) - see documentation to cdd(http://www.ifor.math.ethz.ch/ fukuda/cddman/cddman.html)

Description: This function converts a given configuration (n particles, m Measurements) or a given list of vertices from V-representation into a H-representation. Like above 3.4 the potentially long calculation time has to be taken into consideration, depending on the complexity of the problem.

Example:

```
In[]= ConvToHRep [ 2,3,"2_3"]
Out[]= {{H-representation}, {begin}, {684, 16, real}, {2,
      0,...}..., {end}},
      wheras in this case the files "2_3.ext" (created by Mathematica containing the data for cdd) and "2_3.ine" (created by cdd as result of the calculation) are generated in the current working directory.
```

3.7 ReadInHRep

ReadInHRep [file]

file (String): ".ine" - file containing the H-representation which is to be read in.

Description: Reads the H-representation from a given ".ine"-file for further use in various functions like GetViolInequalities[...] or PlotInequalities[...]

Example:

```
In[]= ReadInHRep[ "2_3"]
Out[]= {{H-representation},{begin},{684, 16,real},{2, 0,...}...,{end}}
```

3.8 GetInequFromHRep

```
GetInequFromHRep[hrep]
```

hrep (List): H-representation yielded for example from the function Read-InhRep [...] or ConvTohRep [...].

Description: Returns the inequalities from a given H-representation as a list. To make the list more readable you can apply InequToRead [...] on it.

Example:

```
In[]= GetInequFromHRep[ConvToHRep[2,1]]
Out[]= {{a1 - alb1 + b1, 1}, {-a1 + alb1, 0}, {alb1 - b1, 0}, {-alb1, 0}}
```

3.9 InequToRead

InequToRead [inequalities]

inequalites (List): List of inequalities yielded from GetInequFromHRep [...]

Description: Makes the list of inequalities yielded from GetInequFromHRep[...] more readable.

Example:

```
In[]= InequToRead [GetInequFromHRep [ConvToHRep [2,1]]] Out[]= a1-a1b1+b1 \le 1 -a1+a1b1 \le 0 a1b1-b1 \le 0 -a1b1 \le 0
```

3.10 GetViolInequalities

GetViolInequalities [hrep, angles, functions, inequ-nr, violation] or GetViolInequalities [file, angles, functions, inequ-nr, violation, opt: options]

hrep (List): H-representation yielded for example from the function Read-Inhrep [...] or ConvTohrep [...].

file (String): File containing the demanded H-representation in cdd format.

- angles (List): List of measurement angles for each particle, whereas the dimension of the list must represent the configuration. If you choose the configuration "2 particles 2 measurements" the list must have the dimension [2,2], in this case the particles a and b are measured each along two axis given by the angles a_1, a_2, b_1 and b_2 , so this parameter has the form $\{\{a_1, a_2\}, \{b_1, b_2\}\}$.
- functions (Symbol): Functions to calculate the quantum mechanical probability of the events. Considering for example two spin- $\frac{1}{2}$ particles in a singlett state, the probability to find the particles both either in spin "up" or both in spin "down" states is given by $\mathbf{P}^{\uparrow\uparrow}(\theta,\phi) = \mathbf{P}^{\downarrow\downarrow}(\theta,\phi) = \frac{1}{2}\sin^2[\frac{(\theta-\phi)}{2}]$, where θ and ϕ are the measurement angles of the particles.

In defining these functions you have to notice, that for all possible events (single events, two-particle-events,...) an appropriate function definition has to exist, each taking a list as parameter (e.g. $Prob[\{x_-, y_-\}] = \frac{1}{2}Sin[(x-y)/2]^2$).

- inequ-nr (List): Range of rows in H-representation used for checking violated inequalities. Specifying this can be useful, if many inequalities have to be evaluated. The form of the parameter is $\{\min, \max\}$ oder All.
- violation (Real): Only inequalities are printed out, that are violated more than this parameter. Default value is 0.
- opt: options: Options can be PrintOut → True/False, which specifies, if the inequalities are printed out during evaluation or not.

Description: GetViolInequalities [...] calculates the discrepancy of inequalities using the given probability functions and therefore the quantum mechanical violation for a distinct adjustment (i. e. special angles) of the detectors and returns all violated inequalities.

Example:

In[] = GetViolInequalities [ConvToHRep [2,2], $\{\{-\frac{\pi}{6},0\},\{0,\frac{\pi}{6}\}\}$, Prob]

Out[]= $(-a1b1+a1b2+a2-a2b1-a2b2+b1 \le 1 \frac{9}{8})$

The probability for the single event of measuring one particle in spin "up" or spin "down" at any angle is given by $Prob[\{x_-\}] = \frac{1}{2}$ and the probability of the joint event has been calculated by $Prob[\{x_-, y_-\}] = \frac{1}{2}Sin[(x-y)/2]^2$.

3.11 PlotInequalities

PlotInequalities [hrep, range, angles, functions, opt: options] or PlotInequalities [file, range, angles, functions, inequ-nr, violation, opt: options]

hrep (List): H-representation yielded for example from the function Read-Inhrep [...] or ConvTohrep [...].

file (String): File containing the demanded H-representation in cdd format.

range (List): Parameter specifying the range for the variable x, which is plotted on the horizontal axis. It has the form $\{x, x_{min}, x_{max}\}$ (see *Mathematica* function Plot [...])

angles (List): List of measurement angles for each particle, whereas the dimension of the list must represent the configuration. If you choose the configuration "2 particles - 2 measurements" the list must have the dimension [2,2], in this case the particles a and b are measured each along two axis given by the angles a_1, a_2, b_1 and b_2 , so this parameter has the form $\{\{a_1, a_2\}, \{b_1, b_2\}\}$.

functions (Symbol): Functions to calculate the quantum mechanical probability of the events. Considering for example two spin- $\frac{1}{2}$ particles in a singlett state, the probability to find the particles both either in spin "up" or both in spin "down" states is given by $\mathbf{P}^{\uparrow\uparrow}(\theta,\phi) = \mathbf{P}^{\downarrow\downarrow}(\theta,\phi) = \frac{1}{2}\sin^2[\frac{(\theta-\phi)}{2}]$, where θ and ϕ are the measurement angles of the particles.

In defining these functions you have to notice, that for all possible events (single events, two-particle-events,...) an appropriate function definition has to exist, each taking a list as parameter (e.g. $Prob[\{x_-, y_-\}] = \frac{1}{2}Sin[(x-y)/2]^2$).

inequ-nr (List): Range of rows in H-representation used for checking violated inequalities. Specifying this can be useful, if many inequalities have to be evaluated. The form of the parameter is {min, max} oder All.

violation (Real): Only inequalities are printed out, that are violated more than this parameter. Default value is 0.

opt: options: Options for the Mathematica- function Plot [...] can be handed over.

Description: This function yields a plot of (violated) inequalities, whereas the function plotted is f(x) = p(x) - b derived from the inequalities of the form $p(x) \le b$ (p(x) is a linear combination of functions to calculate the probabilities of (joint) events, dependent on one variable). Consequently the degree of violation is represented as a positive value of f(x).

Take for example the case "2 particles - 2 measurement directions", where the inequality

$$-p_{a1b1} + p_{a1b2} + p_{a2} - p_{a2b1} - p_{a2b2} + p_{b1} \le 1$$

appears. The probability for the single event of measuring one particle in spin "up" or spin "down" at any angle is given by

$$p_{a1}(x) = p_{a2}(x) = p_{b1}(x) = p_{b2}(x) = \frac{1}{2}$$

and the probability of the joint event can be calculated by

$$p_{a1b1}(x,y) = p_{a2b1}(x,y) = p_{a1b1}(x,y) = p_{a2b2}(x,y) = \frac{1}{2}\sin[(x-y)/2]^2.$$

If we define the measurement angles by

$$a_1 = -\frac{\pi}{3} + x$$
 $a_2 = b_1 = 0$ $b_2 = 2 * \pi$

the inequality can be written as

$$1 + \frac{1}{2}\sin\frac{1}{2}(-\frac{\pi}{3} - x) - \frac{\sin x}{2} - \frac{1}{2}\sin\frac{1}{2}(-\frac{\pi}{3} + x) \le 1$$

The left side is dependent on x $(p(x) = 1 + \frac{1}{2}\sin\frac{1}{2}(-\frac{\pi}{3}-x) - \frac{\sin x}{2} - \frac{1}{2}\sin\frac{1}{2}(-\frac{\pi}{3}+x))$ and b=1. The function to be plotted is f(x)=p(x)-b, thus

$$f(x) = 1 + \frac{1}{2}\sin\frac{1}{2}(-\frac{\pi}{3} - x) - \frac{\sin x}{2} - \frac{1}{2}\sin\frac{1}{2}(-\frac{\pi}{3} + x) - 1$$

Example:

In[]= PlotInequalities [ConvToHRep [2,2], $\{x,0,\pi\}$, $\{\{-\frac{\pi}{3}+x,0\},\{0,2x\}\}$,Prob]

Out[]=

The functions to calculate the probabilities have been defined as $Prob[\{x_-\}] := \frac{1}{2}$ for a single event and $Prob[\{x_-, y_-\}] := \frac{1}{2}Sin[(x-y)/2]^2$ for two-particle events.

3.12 ContPlotInequalities

ContPlotInequalities [hrep, rangex, rangey, angles, func, ineq-nr, violation, opt: options] or

ContPlotInequalities [file, rangex, rangey, angles, func, ineq-nr, violation, opt: options]

 $hrep \ (List): \ H-representation \ yielded \ for \ example \ from \ the \ function \ Read-Inhrep [\dots] \ or \ ConvTohRep [\dots] .$

- rangex (*List*): Parameter specifying the range for the variable x, which is plotted on the horizontal axis. It has the form $\{x, x_{min}, x_{max}\}$ (see *Mathematica* function Plot [...])
- rangex (List): Parameter specifying the range for the variable y, which is plotted on the vertical axis. It has the form $\{y, y_{min}, y_{max}\}$ (see *Mathematica* function Plot [...])
- angles (List): List of measurement angles for each particle, whereas the dimension of the list must represent the configuration. If you choose the configuration "2 particles 2 measurements" the list must have the dimension [2,2], in this case the particles a and b are measured each along two axis given by the angles a_1, a_2, b_1 and b_2 , so this parameter has the form $\{\{a_1, a_2\}, \{b_1, b_2\}\}$.
- functions (Symbol): Functions to calculate the quantum mechanical probability of the events. Considering for example two spin- $\frac{1}{2}$ particles in a singlett state, the probability to find the particles both either in spin "up" or both in spin "down" states is given by $\mathbf{P}^{\uparrow\uparrow}(\theta,\varphi) = \mathbf{P}^{\downarrow\downarrow}(\theta,\varphi) = \frac{1}{2}\sin^2[\frac{(\theta-\varphi)}{2}]$, where θ and φ are the measurement angles of the particles.
 - In defining these functions you have to notice, that for all possible events (single events, two-particle-events,...) an appropriate function definition has to exist, each taking a list as parameter (e.g. $Prob[\{x_-, y_-\}] = \frac{1}{2}Sin[(x-y)/2]^2$).
- inequ-nr (List): Range of rows in H-representation used for checking violated inequalities. Specifying this can be useful, if many inequalities have to be evaluated. The form of the parameter is {min, max} oder All.
- violation (Real): Only inequalities are printed out, that are violated more than this parameter. Default value is 0.

opt: options: Options for the Mathematica-function Plot [...] can be handed over.

Description: Like the function PlotInequalities [...] Cont-PlotInequalities [...] yields a graphical representation of the violation of Boole-Bell type inequalities, but in this case the functions are dependant on two variables: f(x,y) = p(x,y) - b derived from $p(x,y) \le b$, where p(x,y) is a linear combination of functions to calculate the probability for single or joint events. Like in the description of the PlotInequalities [...] -function we take the "2 particles - 2 measurement directions", the only difference is the selection of the measurement angles:

$$a_1 = x$$
 $a_2 = b_1 = 0$ $b_2 = y$

Thus the function f(x, y) is now given by

$$f(x,y) = 1 - \frac{1}{2}\sin\frac{x}{2} + \frac{1}{2}\sin\frac{x-y}{2} - \frac{1}{2}\sin\frac{y}{2} - 1$$

and can be plotted as contour plots. A higher level of violation is represented by a darker contour layer.

Example:

```
\label{eq:local_local_local} \footnotesize \texttt{In[]= ContPlotInequalities[ConvToHRep[2,2],} \{x,0,\pi\}, \{y,0,\pi\}, \{\{x,0\},\{0,y\}\}, \texttt{Prob,All]} \\ \textit{Out[]=}
```

The functions to calculate the probabilities have been defined as $Prob[\{x_-\}] := \frac{1}{2}$ for a single event and $Prob[\{x_-, y_-\}] := \frac{1}{2}Sin[(x-y)/2]^2$ for two-particle events.

3.13 Cdd

Cdd [values, opt: command]

values (List): Data handed over to cdd in an input-file. Each element of the list must be a string and represents a row in the output-file.

opt: command (String): Command to be executed on the generated output-file (default is "cdd" - the default value can be changed utilizing the function SetCd-dCmd [newcommand_String]). Here you can specify for example "cddf" or "cddr".

Description: Simple Interface to run cdd in *Mathematica*. The current version cannot distinguish, whether cdd has produced correct output or not, so you have to pay attention while using this function.

Example:

```
In[] = Cdd [ {"H-Representation", "begin", "6 4 real", "2 -1 0 0", "2 0 -1 0", "-1 1 0 0", "-1 0 1 0", "-1 0 0 1", "4 -1 -1 0", "end"}]
```

```
Out[]= {{"*", "cdd:", "Double", "Description", "Method", "C-Code: Version", "0.61b", "(November", 29, "1997)"}, {"*", "Copyright", "(C)", 1996, "Komei", "Fukuda,", "fukuda@ifor.math.ethz.ch"}, {"*Input", "File:tmp.ine", "(", 6, "x", "4)"}, {"*HyperplaneOrder:", "LexMin"}, {"*Degeneracy", "preknowledge", "for", "computation:", "None", "(possible", "degeneracy)"}, {"*Vertex/Ray", "enumeration", "is", "chosen."}, {"*Computation", "completed", "at", "Iteration", 6.}, {"*Computation", "starts", "at", "Thu", "Mar", 22, "18:48:36", 2001}, {"*", "terminates", "at", "Thu", "Mar", 22, "18:48:36", 2001}, {"*Total", "processor", "time", "=", 0, "seconds"}, {"*", "=", 0, "hour", 0, "min", 0, "sec"}, {"*FINAL", "RESULT:"}, {"*Number", "of", "Vertices", "=", 4, "Rays", "=", 1}, {"V-representation"}, {"begin"}, {5, 4, "real"}, {1, 2, 1, 1}, {1, 1, 1, 1}, {1, 1, 2, 1}, {1, 2, 2, 1}, {0, 0, 0, 1}, {"end"}, {"hull"}}
```

4 Examples

These two examples and are originally solved in a paper by Pitowsky and Svozil [15]. The associated *Mathematica* - notebooks are "3_2.nb" (three particles - 2 measurement directions) and "2_3.nb" (two particles - three measurement directions)

4.1 Three particles and two measurement directions

In this configuration three particles (a, b and c) are measured in detectors which can be switched between two angles each. Consequently there are six different propositions for single particle events: a_1 , a_2 , b_1 , b_2 , c_1 , c_2 , supposing that a_1 is the detection (i. e. the click in a counter) of the particle a in the detector set along the axis specified by the first angle for particle a, b_2 the detection of particle b at the second angle for this particle, and so on ...(cf. Figure 3). If we also take two and three particle events into

Figure 3: Setting for "2 particles - 3 angles"

account (for example the event a_2c_1 means a click in the counter for particle a at the second angle AND a click in the counter for particle c at the first angle), there are 26 different events:

 $a_1, a_2, b_1, b_2, c_1, c_2, a_1b_1, a_1b_2, a_2b_1, a_2b_2, a_1c_1, a_1c_2, a_2c_1, a_2c_2, b_1c_1, b_1c_2, b_2c_1, b_2, c_2, a_1b_1c_1, a_1b_1c_2, a_1b_2c_1, a_1b_2c_2, a_2b_1c_1, a_2b_1c_2, a_2b_2c_1, a_2b_2c_2$

4.2 Violations of inequalities

The truth table for this configuration can be obtained utilizing the function <code>TruthTable[3,2]</code> ³, executing <code>ConvToHRep[3,2]</code> yielded the appropriate H-representation, but this would last quite long, due to the complexity of the correlation polytope for this setting (there are 53856 hyper-planes limiting the polytope). Because of this fact trying to read in the H-representation created by <code>cdd</code> (using <code>ReadInHRep[...]</code>) could also result in memory resource problems by <code>Mathematica</code>.

To avoid this symptoms it is suggested to export the list of vertices and apply cdd outside of *Mathematica* to the file containing the list of vertices (V-representation). This can be done by invoking ToCddExtFile ["3_2",3,2], which creates a file "3_2.ext". This file can be handed over to cdd as parameter to get the file "3_2.ine" comprising the H-representation of the correlation polytope (Command: "cdd 3_2.ext").

Now the search for violated inequalities can begin using the function GetViolinequalities[...]:

May be accepted that the functions to calculate the quantum probabilities of the (joint) events (Prob) have been defined by $Prob[\{x_-\}] := \frac{1}{2}$, $Prob[\{x_-, y_-\}] := \frac{1}{4}$ and $Prob[\{x_-, y_-, z_-\}] := \frac{1}{8} * (1 - Sin[x + y + z])$, where x, y and z are the angles used for detection of each particle,

³due to lack of space not listed here, but it can be found in the *Mathematica*- notebook 3_2.nb

```
GetViolInequalities ["3_2.ine", \{\{0,\frac{\pi}{2}\},\{0,\frac{\pi}{2}\},\{0,\frac{\pi}{2}\}\}\}, Prob, All, 0.4]
yields:
              {{-3 a1+2 a1b1+a1b1c1-4 a1b1c2+3 a1b2-3 a1b2c1-4 a1b1c2+3 a1b2c1-4 a1b1c2+3 a1b2c1-4 a1b1c2+3 a1b2c3-4 a1b2c3-4 a1b2c3-4 a1b1c2+3 a1b2c3-4 a1b2c3-4 a1b2c3-4 a1b1c2+3 a1b2c3-4 a1b2c3
                                    a1b2c2+a1c1+3 a1c2+2 a2b1-2 a2b1c1-a2b1c2-
                                    2 a2b2+a2b2c1+3 a2b2c2+a2c1-a2c2-2 b1+
                                    b1c1+2b1c2+b2c1-2b2c2-c1 \le 0,0.5,
                    {-2 a1+2 a1b1+a1b1c1-4 a1b1c2+2 a1b2-2 a1b2c1-
                                    a1b2c2+a1c1+2 a1c2+3 a2b1-3 a2b1c1-a2b1c2-
                                    2 a2b2+a2b2c1+3 a2b2c2+a2c1-2 a2c2-3 b1+
                                   b1c1+3 b1c2+b2c1-b2c2-c1 \le 0,0.5,
                    {-2 a1+a1b1+a1b1c1-4 a1b1c2+2 a1b2-3 a1b2c1-
                                    a1b2c2+2 a1c1+2 a1c2+2 a2b1-3 a2b1c1-
                                    a2b1c2-a2b2+a2b2c1+2 a2b2c2+a2c1-a2c2-
                                    2 b1+2 b1c1+2 b1c2+b2c1-b2c2-2 c1 \le 0,0.5,
                    {-2 a1+2 a1b1+a1b1c1-4 a1b1c2+2 a1b2-3 a1b2c1-
                                    a1b2c2+3 a1c2+2 a2b1-3 a2b1c1-a2b1c2-
                                    2 a2b2+a2b2c1+2 a2b2c2+a2c1-2 b1+2 b1c1+
                                    2 b1c2+b2c1-b2c2-c1-c2 \le 0,0.5,
                    {-a1+a1b1+a1b1c1-4 a1b1c2+a1b2-3 a1b2c1-
                                    a1b2c2+a1c1+2 a1c2+a2+a2b1-2 a2b1c1-
                                    a2b1c2-2 a2b2+a2b2c1+3 a2b2c2-a2c2-b1+
                                    b1c1+2b1c2+b2+b2c1-2b2c2-c1 \le 1,1.5,
                     {-a1+2 a1b1-a1b1c1-3 a1b1c2+a1b2-4 a1b2c1+
                                    a1b2c2+3 a1c1+a2+a2b1-2 a2b1c1-a2b1c2-
                                    2 a2b2+a2b2c1+3 a2b2c2-a2c2-2 b1+2 b1c1+
                                    2 b1c2+b2+2 b2c1-3 b2c2-2 c1+c2 \le 1,1.5,
                    {-2 a1+2 a1b1-a1b1c1-3 a1b1c2+a1b2-2 a1b2c1-
                                         a1b2c2+2 a1c1+2 a1c2+a2+a2b1-4 a2b1c1+
                                         a2b1c2-2 a2b2+a2b2c1+3 a2b2c2+2 a2c1-3 a2c2-
                                         b1+3 b1c1+b2-b2c2-2 c1+c2 \le 1,1.5
```

4.3 Graphical representation

{...}...}

Using PlotInequalities [...] a graph can be created showing the violation of inequalities dependent on one variable. Defining the probability functions as above, executing

```
PlotInequalities [ "3_2.ine",\{x,0,\pi\},\{\{0,x\},\{0,x\},\{0,x\},Prob,\{10000,20000\},0.4 ]
```

yields the following plot (cf. Figure 4), whereas the corresponding H-representation has to be stored in the file "3_2.ine", {10000,20000} indicates the range of row numbers taken for calculating the graph and 0.4 is the minimal degree of violation to include the inequality in the graph:

```
Figure 4: PlotInequalities["3-2.ine",\{x,0,\pi\},\{\{0,x\},\{0,x\},\{0,x\},Prob,\{10000,20000\},0.4\}
```

To display inequalities dependent on two variables the function Cont-PlotInequalities[...] is provided. This function shows the violation as a contour plot, a more violated set of detection angles results in a darker region in the plot.

```
ContPlotInequalities [ "3.2.ine",\{x,0,\pi\},\{x,0,\pi\},\{\{0,x\},\{0,y\},\{x,y\}\},Prob,\{10000,20000\},0.4 ] returns ContourGraphics-objects, which can be displayed for example by executing
```

```
Show [GraphicsArray [Partition [ cont, 3 [ [ \{1,2\}, A11 ] ] ] ] // TableForm
```

(see figure 5).

Figure 5: ContPlotInequalities["3_2.ine", $\{x,0,\pi\},\{x,0,\pi\},\{\{0,x\},\{0,y\},\{x,y\}\},Prob,\{10000,20000\},0.4]$ Show[GraphicsArray[Partition[cont,3][[$\{1,2\},All$]] // TableForm

4.4 Two particles and three measurement directions

In the case of two particles (a and b) with three properties (whereas the properties are three different angles of the detectors for each particle denoted by a_1 , a_2 , a_3 , b_1 , b_2 , b_3 - see figure 6) 15 different events can be found:

$${a_1, a_2, a_3, b_1, b_2, b_3, c_1, c_2, c_3, a_1b_1, a_1b_2, a_1b_3, a_2b_1, a_2b_2, a_2b_3, a_3b_1, a_3b_2, a_3b_3}$$

Figure 6: Setting for "2 particles - 3 angles"

4.5 Violations of inequalities

Using TruthTable [2,3] all vertices of the corresponding correlation polytope can be found - we get a dimension of 15 and 64 vertices as result (table 2).

a1	a2	a3	b1	b2	b3	alb1	alb2	alb3	a2b1	a2b2	a2b3	a3b1	a3b2	a3b3
0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
1 0	0 1	0	0	0	0	0	0	0	0	0	0	0	0	0
1	1	0	0	0	0	0	0	0	0	0	0	0	0	0
0	0	1	0	0	0	0	0	0	0	ő	0	0	0	ő
1	0	1	0	0	0	0	0	0	0	0	0	0	0	0
0	1	1	0	0	0	0	0	0	0	0	0	0	0	0
1	1	1	0	0	0	0	0	0	0	0	0	0	0	0
0	0	0	1	0	0	0	0	0	0	0	0	0	0	0
1	0	0	1	0	0	1	0	0	0	0	0	0	0	0
0 1	1 1	0	1 1	0	0	0 1	0	0	1 1	0	0	0	0	0
0	0	1	1	0	0	0	0	0	0	0	0	1	0	0
1	0	1	1	0	0	1	0	0	0	Ö	0	1	0	Ö
0	1	1	1	0	0	0	0	0	1	0	0	1	0	0
1	1	1	1	0	0	1	0	0	1	0	0	1	0	0
0	0	0	0	1	0	0	0	0	0	0	0	0	0	0
1	0	0	0	1	0	0	1	0	0	0	0	0	0	0
0	1	0	0	1	0	0	0	0	0	1	0	0	0	0
1	1	0	0	1	0	0	1	0	0	1	0	0	0	0
0 1	0	1 1	0	1 1	0	0	0 1	0	0	0	0	0	1 1	0
0	1	1	0	1	0	0	0	0	0	1	0	0	1	0
1	1	1	0	1	0	0	1	ő	0	1	0	0	1	0
0	0	0	1	1	0	0	0	0	0	0	0	0	0	0
1	0	0	1	1	0	1	1	0	0	0	0	0	0	0
0	1	0	1	1	0	0	0	0	1	1	0	0	0	0
1	1	0	1	1	0	1	1	0	1	1	0	0	0	0
0	0	1	1	1	0	0	0	0	0	0	0	1	1	0
1 0	0 1	1 1	1 1	1 1	0	1	1	0	0 1	0 1	0	1 1	1 1	0
1	1	1	1	1	0	1	1	0	1	1	0	1	1	0
0	0	0	0	0	1	0	0	0	0	0	0	0	0	0
1	Ö	0	Ö	0	1	ő	0	1	Ö	Ö	0	ő	ő	Ö
0	1	0	0	0	1	0	0	0	0	0	1	0	0	0
1	1	0	0	0	1	0	0	1	0	0	1	0	0	0
0	0	1	0	0	1	0	0	0	0	0	0	0	0	1
1	0	1	0	0	1	0	0	1	0	0	0	0	0	1
0 1	1 1	1 1	0	0	1 1	0	0	0 1	0	0	1 1	0	0	1 1
0	0	0	1	0	1	0	0	0	0	0	0	0	0	0
1	0	0	1	0	1	1	0	1	0	0	0	0	0	0
0	1	0	1	0	1	0	0	0	1	0	1	0	0	0
1	1	0	1	0	1	1	0	1	1	0	1	0	0	0
0	0	1	1	0	1	0	0	0	0	0	0	1	0	1
1	0	1	1	0	1	1	0	1	0	0	0	1	0	1
0	1	1	1	0	1	0	0	0	1	0	1	1	0	1
1	1	1	1	0	1	1	0	1	1	0	1	1	0	1
0 1	0	0	0	1 1	1 1	0	0 1	0 1	0	0	0	0	0	0
0	1	0	0	1	1	0	0	0	0	1	1	0	0	0
1	1	0	0	1	1	0	1	1	0	1	1	0	0	0
0	0	1	0	1	1	0	0	0	0	0	0	0	1	1
1	0	1	0	1	1	0	1	1	0	0	0	0	1	1
0	1	1	0	1	1	0	0	0	0	1	1	0	1	1
1	1	1	0	1	1	0	1	1	0	1	1	0	1	1
0	0	0	1	1	1	0	0	0	0	0	0	0	0	0
1	0	0	1	1	1	1	1	1	0	0	0	0	0	0
0 1	1 1	0	1 1	1 1	1 1	0 1	0 1	0 1	1 1	1 1	1 1	0	0	0
0	0	1	1	1	1	0	0	0	0	0	0	1	1	1
0	-	•	•	•	-	-		-	-	•	9	•	-	-

a1	a2	a3	b1	b2	b3	alb1	alb2	alb3	a2b1	a2b2	a2b3	a3b1	a3b2	a3b3
1	0	1	1	1	1	1	1	1	0	0	0	1	1	1
0	1	1	1	1	1	0	0	0	1	1	1	1	1	1
1	1	1	1	1	1	1	1	1	1	1	1	1	1	1

Table 2: Truth table for 6 propositions

Executing hrep=ConvToHRep [2,3] a H-representation of the polytope described by the truth table above can be created. This results in 684 hyper-planes respectively 684 inequalities from the 64 vertices limiting the polytope:

```
{{H-representation}, {begin},{684, 16, real}, {2, 0, -2, 1, -1, 0, -1, 1, -1, 0, 1, 1, 1, -1, -1, 1}, {2, -2, 0, 1, -1, 0, -1, 1, 1, 1, 1, -1, 0, -1, -1, 1}, {...}, {1, -1, 0, 0, -1, 0, 0, 1, 0, 0, 0, 0, 0, 0, 0}, {end},{Konfiguration,2,3}}
```

All inequalities can be displayed by

```
GetInequFromHRep[hrep]// InequToRead
```

The result of this operation is

```
-a1b1 + a1b2 + 2 a2 - a2b1 - a2b2 - a2b3 - a3 + a3b1 + a3b2 - a3b3 + b1 + b3 \leq 2 2a1 - a1b1 - a1b2 - a1b3 - a2b1 + a2b2 - a3 + a3b1 + a3b2 - a3b3 + b1 + b3 \leq 2 a1 - a1b1 - a1b3 - a2b1 + a2b2 + a2b3 - a3 + a3b1 + a3b2 - a3b3 + b1 - b2 \leq 1 \ldots a1 - a1b1 + b1 \leq 1
```

Using

```
GetViolInequalities [hrep, \{\{0,\frac{2\pi}{3},\frac{4\pi}{3}\},\{\{0,\frac{2\pi}{3},\frac{4\pi}{3}\}\},Prob,All] // TableForm
```

all inequalities can be displayed that are violated at the specific angles $a_1=b_1=0$, $a_2=b_2=\frac{2\pi}{3}$ and $a_3=b_3=\frac{4\pi}{3}$ taking the functions $\text{Prob}[\{x_-\}]:=\frac{1}{2}$ and $\text{Prob}[\{x_-\}]:=\sin[\frac{x-y}{2}]/2$ to calculate the quantum probabilities, which is equivalent to the probability to find two spin- $\frac{1}{2}$ particles in a singlet state $(|\psi\rangle=\frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle)-|\downarrow\uparrow\rangle)$ both either in spin "up" or both in spin "down".

4.6 Graphical representation

Like in the configuration "three particles - two angles" described above a graphical representation can be generated either dependent on one or dependent on two variables. On the one hand side by using

PlotInequalities [hrep,
$$\{x,0,\pi\}$$
, $\{\{0,x,\frac{4\pi}{3}\},\{\{0,x,\frac{4\pi}{3}\}\}$,Prob] we get a plot of all violated inequalities (cf. Figure 7),

Figure 7: PlotInequalities[hrep,
$$\{x,0,\pi\}$$
, $\{\{0,x,\frac{4\pi}{3}\},\{\{0,x,\frac{4\pi}{3}\}\}$,Prob]

on the other hand side contour plots of all inequalities violated for example more than 0.2 can be generated by executing

cont=ContPlotInequalities[hrep,
$$\{x,0,\pi\}$$
, $\{y,0,\pi\}$, $\{\{0,x,y\}$, $\{\{0,x,y\}\}$,Prob,All,0.2]

To display the outcome of the calculation entering Show /@ cont results in ContourPlots of the following form (cf. Figure 8):

Figure 8: ContPlotInequalities[hrep, $\{x,0,\pi\}$, $\{y,0,\pi\}$, $\{\{0,x,y\}\}$,Prob,All,0.2] Show /@ %

References

- [1] Itamar Pitowsky. Quantum Probability—Quantum Logic. Springer, Berlin, 1989.
- [2] Itamar Pitowsky. Correlation polytopes their geometry and complexity. *Mathematical Programming*, 50:395–414, 1991.
- [3] George Boole. On the theory of probabilities. *Philosophical Transactions of the Royal Society of London*, 152:225–252, 1862.
- [4] Komei Fukuda. cdd program. 2000.
- [5] George Boole. *An investigation of the laws of thought*. Dover edition, New York, 1958.
- [6] Theodore Hailperin. *Boole's logic and probability (Studies in logic and the foundations of mathematics ; 85).* North-Holland, Amsterdam, 1976.
- [7] Itamar Pitowsky. George Boole's 'conditions od possible experience' and the quantum puzzle. *Brit. J. Phil. Sci.*, 45:95–125, 1994.
- [8] J. F. Clauser and M. A. Horne. *Physical Review*, D10:526, 1974.
- [9] J. F. Clauser, M. A. Horne, A. Shimony, and R. A. Holt. *Physical Review Letters*, 23:880–884, 1969.
- [10] J. F. Clauser and A. Shimony. Bellos theorem: experimental tests and implications. *Rep. Prog. Phys.*, 41:1881–1926, 1978.
- [11] Itamar Pitowsky. From George Boole to John Bell: The origin of Bell's inequality. In M. Kafatos, editor, *Bell's Theorem, Quantum Theory and the Conceptions of the Universe*. Kluwer, Dordrecht, 1989.
- [12] Günter M. Ziegler. Lectures on Polytopes. Springer, New York, 1994.
- [13] T.S. Motzkin, H. Raiffa, G.L. Thompson, and R.M. Thrall. The double description method. In *Contributions to theory of games, Vol.* 2. Princeton University Press, New Jersey, Princeton, 1953.
- [14] A. N. Kolmogorov. *Grundbegriffe der Wahrscheinlichkeitsrechnung*. Springer, Berlin, 1933. English translation in [16].
- [15] Itamar Pitowsky and Karl Svozil. New optimal tests of quantum nonlocality. *Physical Review*, A, to appear, 2001. arXiv:quant-ph/0011060.
- [16] A. N. Kolmogorov. *Foundations of the Theory of Probability*. Chelsea, New York, 1956.